

High Pressure Research



An International Journal

ISSN: 0895-7959 (Print) 1477-2299 (Online) Journal homepage: http://www.tandfonline.com/loi/ghpr20

Raman analysis of an impacted α -GeO₂–H₂O mixture

Ivonne Rosales , Claude Thions-Renero , Erendira Martinez , Fernando Agulló-Rueda , Lauro Bucio & Eligio Orozco

To cite this article: Ivonne Rosales , Claude Thions-Renero , Erendira Martinez , Fernando Agulló-Rueda , Lauro Bucio & Eligio Orozco (2012) Raman analysis of an impacted α -GeO $_2$ -H $_2$ O mixture, High Pressure Research, 32:3, 396-404, DOI: 10.1080/08957959.2012.710233

To link to this article: http://dx.doi.org/10.1080/08957959.2012.710233



Full Terms & Conditions of access and use can be found at http://www.tandfonline.com/action/journalInformation?journalCode=ghpr20



Raman analysis of an impacted α-GeO₂–H₂O mixture

Ivonne Rosales^a, Claude Thions-Renero^b, Erendira Martinez^a, Fernando Agulló-Rueda^c, Lauro Bucio^a* and Eligio Orozco^a

^aDepartamento de Estado Sólido, Instituto de Física, Universidad Nacional Autónoma de México, México City, México; ^bDepartamento de Física Teórica, Instituto de Física, Universidad Nacional Autónoma de México, México City, México; ^cInstituto de Ciencia de Materiales de Madrid, CSIC, Madrid, Spain

(Received 24 April 2012; final version received 14 June 2012)

Through a Raman analysis, we detected polymorphism at high pressure on mixtures of α -GeO₂ microcrystalline powder and water under impact experiments with a single-stage gas gun. The Raman measurements taken from recovered samples show two vibrational modes associated with water-related species. After the impact, the size of the α -GeO₂ crystallites was approximately 10 times higher showing molten zones and a lot of porous faces. Raman examination showed some unknown peaks possibly associated with other GeO₂ polymorphs detected by X-ray diffraction experiments and perhaps stabilized in the porous of the α -GeO₂ crystallites.

Keywords: phase transformation; shock waves; germanium dioxide; polymorphism

1. Introduction

The study of chemical reactions and phase transformations occurring at high pressures and temperatures is an interesting topic in fields such as earth and planetary sciences as well as materials science. Concerning the first two cases, particular importance has been given to the phenomena of forming minerals which appear during the impact processes of meteors with earth or planets and, also, the physicochemical processes occurring between minerals inside the earth mantle [1]. In the case of materials science, two main motivations are focussed for the research of this kind of transformations: one of them deals with the basic physical and chemical principles of materials working at high pressure/temperatures, while the other is concerned with the rise of new materials having novel physical properties and potential technological applications [2].

In general, from the experimental point of view, two techniques are used to induce chemical reactions and phase transformations at high pressures and temperatures: the static compression method and the shock compression method. In the first method, the pressure is applied to the sample gradually by using devices such as diamond anvil cells, where the hydrostatic pressure can reach as high as 100 GPa. Simultaneously, the sample can be heated in a controlled manner using, e.g. Nd-YAG kind of laser. In these experiments, the detection of the phases that formed during the compression process is done *in situ*, using different techniques such as Raman spectroscopy (RE)

ISSN 0895-7959 print/ISSN 1477-2299 online © 2012 Taylor & Francis

http://dx.doi.org/10.1080/08957959.2012.710233

http://www.tandfonline.com

^{*}Corresponding author. Email: bucio@fisica.unam.mx

and X-ray synchrotron diffraction, which are the most commonly used. The latter allows a more detailed analysis of phase transformations induced by pressure and temperature applied in such a small region corresponding to the small size of the samples, only a few tens of microns [3]. In the shock compression method, high pressure pulses (microseconds in duration) are applied to the sample by using explosives or projectiles traveling at speeds between 1 and 10 km/s and impact the samples encapsulated in metal containers. In the latter case, gas guns of one or two stages have been used; these are loaded with pressurized gases (He and H) that are released abruptly to accelerate the projectiles [4]. Other devices have been recently developed that speed up magnetically metal projectiles that reach velocities somewhat higher than 20 km/s [5]. Whether using explosives or projectile accelerator devices, the shock compression method can induce physicochemical changes on the sample, in time scales at the peak pressure state (shock-induced processes) or can result in the generation of highly activated materials that significantly enhanced their chemical reactivity during post-shock thermal treatments (shock-assisted processes) [6]. To estimate the peak shock pressure and mean temperature on the metal bulk containers, numerical models have been developed for a specific experimental setup in the case of experiments with explosives [7,8]. In experiments with projectiles, the peak shock pressure on the metal container is calculated by applying the impedance match method [4] and the mean bulk temperature is measured directly by spectroscopic techniques [9] or estimated from high pressure/temperature equation of state of impacted materials. This methodology, together with the microstructural analysis of samples recovered after impact, allows having an overview of the conditions under which there are phase transformations and chemical reactions in the materials under study. Although considerable progress has been made to explain from a theoretical framework mechanism that gives rise to phase transformations and chemical reactions for mixtures of powders of geophysical and technological interest [10], it is still necessary to verify whether the models developed can be applied to other systems, as for example porous samples of rare-earth oxides or of elements of group IV.

Recently, our group has initiated a systematic study of phase transformations that occur in oxides porous samples of type XO_2 (X = Ge, Sn, Pb), which were recovered after being subjected to impact processes by the use of one-stage gas gun. These compounds have the advantage that there is enough information in the literature concerning the phases that are generated in them, when they are subjected to high pressures in diamond cells. In particular, the GeO₂ exhibits a large variety of amorphous and liquid phases under pressure [11]. In a previous work, we have reported fusion processes, phase transformations and grain growth in a recovered sample of mixtures of microcrystalline α -GeO₂ powders and water, submitted to high pressure planar shock waves [12]. After impact, 88% of the sample is α -GeO₂, and the remaining (12%) was found to be transformed to other polymorphs. The amounts of the other phases reported were: 6.0% of monoclinic GeO₂ (S.G. P2₁/c); 4.9 % of CaCl₂-type GeO₂ (S.G. Pnnm); and 1.1 % of rutile-type GeO₂ (S.G. P4₂/mnm). This quantitative phase analysis was performed by Rietveld refinement using conventional X-ray powder diffraction measurements on the recovered sample [12]. All these phases have already been reported qualitatively by different authors, but from static pressure experiments [13]. These results indicate that the shock waves cause the sample to reach temperatures as high or higher than the melting point of α -GeO₂ ($T \ge 1389$ K), and pressures at least equal to the transition of rutile-type GeO_2 to $CaCl_2$ -type GeO_2 ($P \ge 26$ GPa).

The aim of the present research is to report additional data on these phase transformations by performing RE, and scanning electron microscopy (SEM) measurements.

2. Methods

A series of shock experiments on mixtures of deionized water and microcrystalline alpha-GeO₂ powders were carried out using the one-stage light gas gun installed at the High Pressure

Laboratory of the Instituto de Física at the Universidad NacionalAutónoma de México [14]. The propellant gas that was employed in all cases was helium. The projectile and target designs are described elsewhere [12]. With this arrangement, it is possible to perform planar shock-wave experiment by impacting a stainless steel plate against a target which consists of three stainless steel capsule chambers hermetically sealed. Each chamber contains a specific amount of microcrystalline powder of α -GeO₂. In order to promote particle interactions during shock-loading conditions, less dense samples were prepared, mixing α-GeO₂ powders (Sigma-Aldrich 99.99%) with different proportions of deionized water. For the experiment, we employed a plate of 106 g with a thickness of 3.46 mm. The plate was mounted on the nose of a nylon projectile (impact 2000) accelerated to a projectile velocity of 829 m/s; under these conditions, by impedance matching, the peak pressure was estimated to be 14 GPa in the target. SEM observations were performed on a JEOL JSM 5600 LV microscope. Raman spectra were measured with a Renishaw Ramascope 2000 spectrometer by employing a 514.5 nm argon laser line, with a spatial resolution of $1-2 \mu m$. Light was focussed on the sample with an optical microscope ($50 \times$ objective), which also collected the scattered light onto the spectrometer. Laser power density on the sample surface was of the order of 0.2 MW/cm². X-ray powder diffraction of the recovered sample was measured at room temperature in a Bruker D8 Advance diffractometer; Cu Ka1 radiation, and goniometer with a LynxEye detector. The data were collected from $2\theta = 7 - 110^{\circ}$, with a step size of 0.019° and operating conditions of 35 kV and 25 mA in the X-ray generator [12].

3. Results and discussion

Figure 1(a,b) shows at the same amplification, SEM images of powders of the α -GeO₂ phase before and after the shock event, respectively. Before impact, the α -GeO₂ phase shows rhombohedrally faceted crystallites of 1 μ m size. In Figure 1(b), the sample appears with molten zones in faceted crystals showing {100} and {101} forms with size around 10 μ m. It seems that this is due to collisions between particles induced by the shock event. From those figures, it is evident that the impact induced the growth of microcrystals of tens of microns. Moreover, a most detailed SEM analysis of the recovered sample shows zones characterized by partially molten-particle aggregates (Figure 2(a)) and zones with well-crystallized grains (Figure 2(b)). This complex process leads to a heterogeneous distribution of particles of different sizes and, as it will be discussed later, to different high pressure phases associated with GeO₂ as well as to chemical reactions between microcrystalline powders of α -GeO₂ and the added water of the mixture. All this is a consequence of localized pressure gradients and elevated temperatures induced on the sample by the shock pressure. This type of transformations is similar to other previously reported for low-density porous materials [15,16]. In order to detect the distribution of phases generated by the shock waves in the sample, a detailed Raman analysis in different zones was performed.

By analyzing different points of the sample before impact, we obtain the characteristic Raman spectrum of the α -GeO₂ phase (Figure 3(a)), where the inset is an image obtained by optical microscopy that corresponds to the analyzed zone. However, when measuring the powder of the sample recovered, significant variations on the shape of the Raman spectra were found (Figure 3(b)). In this figure, four Raman spectra are plotted and the inset image shows a cluster of microcrystals at which these Raman spectra were taken at four different points. Besides the characteristic bands corresponding to the α -GeO₂ phase, some additional peaks appear, and their presence, intensity and position differ in each spectrum. For some faceted crystallites, we obtained spectra with a combination of narrow and broad bands (Figure 3(c)) what was interpreted as due to the presence of a α -GeO₂ crystallite (narrow bands) with amorphous GeO₂ (broad bands) which could have been formed and stabilized in the porous part of the faceted crystallite (inset SEM image) after the change produced by the shock wave. This kind of spectra has been observed

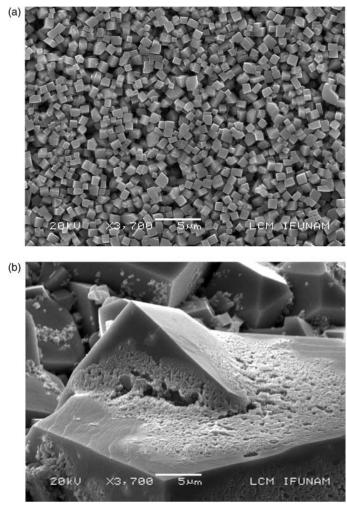
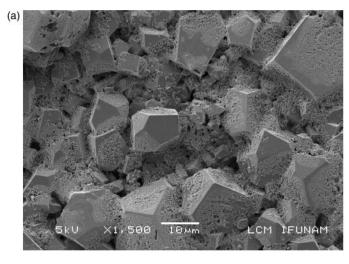


Figure 1. (a) Panoramic SEM image of α -GeO₂ sample before impact. (b) Sample after impact appears with high degree of molten zones with crystallites of dimensions bigger than those from starting α -GeO₂ obtained after impact.

by other researchers on samples of crystalline α -GeO₂ powders submitted to static compression and released in the range from 8 to 20 GPa [17]; and from experiments in which monolithic powders of α -GeO₂ were amorphized by impact [18]. This kind of spectra has been observed by other researchers on samples of crystalline α -GeO₂ powders submitted to static compression and released in the range from 8 to 20 GPa [19] and from experiments in which monolithic powders of α -GeO₂ were amorphized by impact [17]. The Raman spectrum plotted in Figure 4(a) corresponds to the inset image of α -GeO₂ crystal (signaled with an arrow) with well-defined faces. The crystal forms typically observed in α -GeO₂ crystals are represented also in the inset image in which it is possible to appreciate the characteristic rhombohedral symmetry. This Raman spectrum presents well-defined bands which belong to α -GeO₂ with other additional peaks marked with arrows in Figure 4(a). The description of the spectrum is as follows: a double degenerate modes of E symmetry split in transverse optic (TO) and longitudinal optic (LO) modes of α -GeO₂ are localized at the values 121 (TO + LO), 165 (TO + LO), 210 (TO), 325 (TO and a shocking A₁ mode), 517 (LO), 589 (LO) 855 (TO), 962 (TO), 989 (LO) cm⁻¹ and the symmetric shocking A₁ modes at 262, 442, 880 cm⁻¹ being the peak at 442 cm⁻¹, a very intense band that characterizes



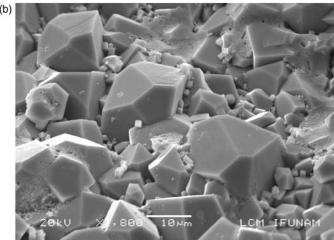


Figure 2. (a) Sample after impact appears with high degree of molten zones similar to those shown in Figure 1(b). (b) Sample after impact appears with low degree of molten zones. Faceted crystals of α -GeO₂ showing typically {100} and {101} forms with size around 10 μ m can be seen.

this polymorph [18,20,21] accompanied by a peak and a shoulder at 370–390 cm⁻¹, assigned to other E symmetry (LO) and (TO) modes of α -GeO₂ by Scott [20]. It is worthy to note that three additional peaks have been detected; a broad band between 700 and 800 cm⁻¹, at higher frequencies, above 1150 cm⁻¹, an unassigned peak to any known GeO₂ phase, and other broad band at 1200–1700 cm⁻¹. A deconvolution of the broad band between 700 and 800 cm⁻¹ (Figure 4(b)) gives two bands; one at 741 and the other at 771 cm⁻¹. According to the Raman active modes for rutile-type GeO₂ given by $\Gamma = A_{1g} + B_{1g} + B_{2g} + E$, the band at 741 cm⁻¹ could be assigned to the active Raman mode A_{1g} [22]. In this case, the change in the Raman shift corresponding to the A_{1g} mode from 701 cm⁻¹ to a higher value could be a consequence of the pressure applied during the impact [18,22]. These and other bands for different GeO₂ polymorphs, which have been characterized by RE [18–22], are summarized in Table 1. The band localized at 771 cm⁻¹, has a position closer to the one reported for a peak assigned to the stretching mode of the Ge–OH complex (769 cm⁻¹) in GeO₂ gels [23]. Since in our experiment, the sample of α -GeO₂ under shock compression was mixed with water, it is reasonable to expect that water species appear in

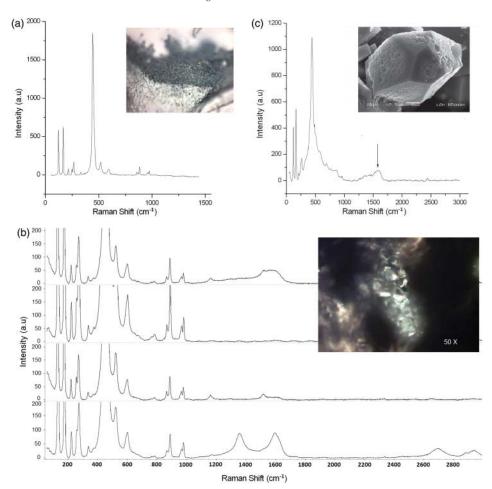


Figure 3. (a) Representative Raman spectrum of the α -GeO₂ obtained by analyzing different points of the sample before impact. Inset figure is an image obtained by optical microscopy that corresponds to the analyzed zone. (b) Raman spectra obtained from the recovered sample. Significant variations on the shape of the bands can be seen. The inset image shows a cluster of microcrystals for which these Raman spectra were taken at four different points. (c) Representative Raman spectrum for some faceted crystallites is presented in the inset image. Narrow and broad bands were interpreted as due to the presence of a α -GeO₂ crystallite (narrow bands) with amorphous GeO₂ (broad bands), which could have been formed and stabilized in the porous part of the faceted crystallite after the change produced by the shock wave.

the Raman spectra. Under shock pressure, it has been proposed that the water molecule transforms according to the following dissociations reactions [24–26]:

$$2H_2O \rightarrow OH^{-1} + H_3O^{+1}$$
 (1)

and

$$H_2O \to OH^{-1} + H^{+1}$$
 (2)

So, the presence of hydroxyl (OH^{-1}) , hydronium (H_3O^{+1}) and protons (H^{+1}) during the impact event, gives the possibility for chemical reactions between microcrystalline α -GeO₂ (or polymorphs) and water. This could be the origin of the hydrous-species encountered in this work in a single micro-crystal (Figure 4(a)). After deconvolution, two bands can be appreciated in Figure 4(c). One of them, centered at 1592 cm⁻¹, could be assigned to the E anti-symmetric vibrational mode of hydronium ion H_3O^{+1} , possibly adsorbed at the surface of α -GeO₂ [27,28],

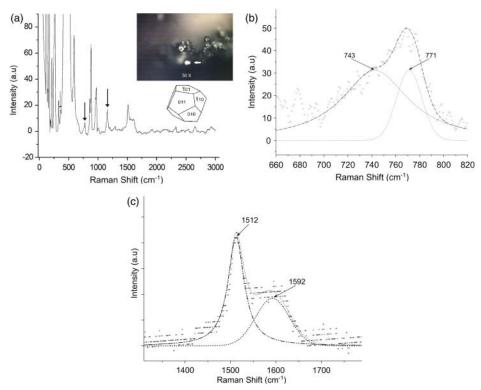


Figure 4. (a) Raman spectrum corresponding to the inset image of α -GeO₂ crystal (signaled with an arrow) with well-defined faces. The crystal forms typically observed in α -GeO₂ crystals are also represented in the inset image in which it is possible to appreciate the characteristic rhombohedral symmetry. (b) The deconvolution of the broad band between 700 and $800\,\mathrm{cm^{-1}}$ gives two bands: one at 741 and the other at 771 cm⁻¹. The band at 741 cm⁻¹ could be assigned to the active Raman mode A_{1g} for rutile-type GeO₂ [22]. (c) Deconvolution of two bands can be appreciated: one of them, centered at $1592\,\mathrm{cm^{-1}}$, could be assigned to the E anti-symmetric vibrational mode of hydronium ion H_3O^{+1} , possibly adsorbed at the surface of α -GeO₂.

Table 1. Active modes of vibration (cm⁻¹) observed in RE for GeO₂ polymorphs. The symmetry mode corresponding for each Raman shift is written below its value.

GeO ₂ polymorph	Raman shift (cm ⁻¹)											
α-Quartz-type GeO ₂ [9–11]	973 E _{LO}	960 E _{TO}	881 A ₁	860 E _{TO}	593 E _{LO}	516 E _{LO}	444 A 1	330 A_{I}, E_{TO}	263 A ₁	212 E _{TO}	166 E ^a	123 E ^a
Rutile-type GeO ₂ (at 0.1 MPa) [12]	LLO	210	211	873 B _{2g}	701 A_{Ig}	LLO	717	117, 270	211	210	L	171 B_{Ig}
CaCl ₂ -type GeO ₂ (at 34.2 GPa) [12]	$_{B_{Ig}}^{1012}$			$815 A_g$	1g							$155 A_g$

 $^{^{}a}TO + LO.$

whereas the other peak, centered at 1512 cm⁻¹, could not be assigned to other vibrational modes of the hydronium or others hydrous-species within the GeO₂ matrix.

The experimental design described in this and in the previous work [12], allows studying the effect of shock waves in three chambers impacted simultaneously, in such a way that the same impact conditions can be ensured on porous oxide samples of different densities. The analysis of the recovered samples, by means of some common techniques such as X-ray powder diffractometry, SEM, micro-RE and others, allows to determine the type of physicochemical changes

induced by the shock compression. In order to detect the pressure and temperatures generated during and/or after the impact on complex systems such as mixtures of reactive compounds, it could be used as temperature and pressure sensors, ceramic oxides as those of group IV, knowing in advance that all of them show the following phase transformation sequence [29]:

```
\alpha-GeO<sub>2</sub> (P3<sub>2</sub>2) \rightarrow rutile-type(P4<sub>2</sub>/mnm at 7 GPa) \rightarrow CaCl<sub>2</sub>-type
(Pnnm at 26 GPa) \rightarrow \alpha-PbO<sub>2</sub>-type (Pbcn at 44 GPa) \rightarrow pyrite-type (Pa-3 at 70 GPa)
```

Following this experimental procedure, it is possible to explore interaction processes between particles of different sizes and/or hardness, in microcrystalline or nano-crystalline, of the same phase or of different phases and of course, chemical reactions between solid–solid, solid–liquid or other kind of reactive powder mixtures as has been pointed out by Bastanov [30]. Our group has started a systematic research work to explore the possible application of this kind of oxides as a high pressure–temperature sensors on impact experiments.

4. Conclusions

The Raman analysis performed in this investigation, allowed us to detect crystalline and crystalline-amorph mixed phases on impacted samples of $\alpha\text{-GeO}_2$. We detected a complex and heterogeneous distribution of those phases on the recovered sample. The Raman spectra of individual micro-crystals (tens of microns) showed all bands associated with longitudinal and traverse optic modes of $\alpha\text{-GeO}_2$ together with additional bands. The bands at 395, 1150 and 1500 cm $^{-1}$ could not be assigned to any known species of GeO $_2$ or other hydrous-species within GeO $_2$. The bands at 471 and at 441 cm $^{-1}$ were assigned as follows: the former to the stretching mode of the GeO-OH in the crystal network and the second to the active Raman mode A_{1g} of rutile GeO $_2$ (400 cm $^{-1}$), shifted to a higher frequency, possibly due to the effect of impact pressure. Our analysis shows that the impact also induces chemical reactions between microcrystalline powders and the water of the mixture, allowing the formation of hydrous species into the amorphous or crystalline structures of GeO $_2$. SEM observations unequivocally show grain growth with partially melted crystalline and well-defined crystalline zones on impacted samples of $\alpha\text{-GeO}_2$.

Acknowledgements

The authors acknowledge the technical support of Cristina Zorrilla and Laboratorio Central de Microscopia, IFUNAM. This work was supported by the projects PAPIIT 105810, CONACYT 81700 and 167624.

References

- Y. Syono and M.H. Manghnani (eds.), High-Pressure Research: Application to Earth and Planetary Sciences, Geophysical Monograph Series Vol. 67, Mineral Physics Vol. 3, 1992, American Geophysical Union, Washington D.C., p. 530.
- [2] H.D. Hochheimer (ed.), Frontiers of High Pressure Research II: Application of High Pressure to Low-Dimensional Novel Electronic Materials, North Atlantic Treaty Organization, Scientific Affairs Division Vol. 2, Kluwer Academic Publishers, The Netherlands, 2001, p. 557.
- [3] W.B. Holzapfel and N. Isaacs (eds.), High Pressure Techniques in Chemistry and Physics, Oxford University Press, Oxford, 1997.
- [4] T. Goto and Y. Syono, Technical aspect of shock compression experiments using the gun method, in Materials Science of the Earth's Interior, I. Sunagawa, ed., Terra Scientific Publishing Company, Tokyo, 1984, pp. 605–619.
- [5] M.D. Knudson, R.W. Lemke, D.B. Hayes, C.A. Hall, C. Deeney, and J.R. Asay, Near-absolute Hugoniot measurements in aluminum to 500 GPa using a magnetically accelerated flyer plate technique, J. Appl. Phys. 94 (2003), pp. 4420–4431.

- [6] K.S. Vandersall and N.N. Thadhani, Investigation of "shock-induced" and "shock-assisted" chemical reactions in Mo⁺²Si powder mixtures, Metall. Mater. Trans. A 34 (2003), pp. 15–23.
- [7] B.O. Garcia and D.J. Chavez, Shock compression of liquid hydrazine, in Proceedings of the Conference of the American Physical Society Topical Group on Shock Compression of Condensed Matter, S.C. Schmidt and W.C. Tao, eds., AIP Conference Proceedings Vol. 370, 1996, pp. 831–834.
- [8] K.S. Vandersall, Investigation of shock-induced and shock-assisted chemical reactions in Mo-Si powder mixtures, Ph.D. thesis, Georgia Institute of Technology, Atlanta, GA, 1999.
- [9] M.B. Boslough, Shock-wave properties and high-pressure equations of state of geophysically important materials, Ph.D. thesis, California Institute of Technology, Pasadena, CA, 1984, p. 171.
- [10] M.A. Meyers, L.-H. Yu, and K.S. Vecchio, Shock synthesis of silicides-II. Thermodynamics and kinetics, Acta Metall. Mater. 42 (1994), pp. 715–729.
- [11] M. Micoulaut, L. Cormier, and G.S. Henderson, *The structure of amorphous, crystalline and liquid GeO*₂, J. Phys., Condens. Matter. 18 (2006), pp. R753–R784.
- [12] I. Rosales, C. Thions-Renero, E. Martinez, L. Bucio, and E. Orozco, Grain growth and phase transformations induced by shock waves on alpha-GeO₂ powder, High Press. Res. 31 (2011), pp. 428–435.
- [13] V.P. Prakapenka, G. Shen, L.S. Dubrovinsky, M.L. Rivers, and S.R. Sutton, *High pressure induced phase transformation of SiO₂ and GeO₂: difference and similarity*, J. Phys. Chem. Solids 65 (2004), pp. 1537–1545.
- [14] M. De Icaza, C.T. Renero, and F. Prieto, Experimental facilities for impact physics research at the National University of Mexico, Phys. B + C 139–140 (1986), pp. 599–602.
- [15] M.B. Boslough and R.A. Graham, Submicrosecond shock-induced chemical reactions in solids: First real-time observations, Chem. Phys. Lett. 121 (4 and 5) (1985), pp. 446–452.
- [16] N.N. Thadhani, R.A. Graham, T. Royal, E. Dunbar, M.U. Anderson, and G.T. Holman, Shock-induced chemical reactions in titanium-silicon powder mixtures of different morphologies: Time-resolved pressure measurements and materials analysis, J. Appl. Phys. 82 (1997), pp. 1113–1128.
- [17] N. Suresh, G. Jyoti, S.C. Gupta, and S.K. Sikka Sangeeta, Shock-induced amorphization in q-GeO₂, J. Appl. Phys. 76 (3) (1994), pp. 1530–1534.
- [18] M. Madon, Ph. Gillet, Ch. Julian, and G.D. Price, A vibrational study of phase transitions among the GeO₂ polymorphs, Phys. Chem. Miner. 18 (1991), pp. 7–18.
- [19] G.H. Wolf, S. Wang, Ch.A. Herbst, D.J. Durben, W.F. Oliver, Z.C. Kang, and K. Holvorson, Pressure induced collapse of the tetrahedral framework in crystalline and amorphous GeO₂, in High-Pressure Research: Application to Earth and Planetary Sciences, Y. Syono and M.H. Manghnani, eds., American Geophysical Union, Washington D.C., 1992, pp. 503–517.
- [20] J.F. Scott, Raman spectra of GeO₂, Phys. Rev. B 1 (1970), pp. 3488–3493.
- [21] W. Dultz, M. Quilichini, J.F. Scott, and G. Lehman, Phonon spectra of quartz isomorphs, Phys. Rev. B 11 (1975), pp. 1648–1653.
- [22] J. Haines, J.M. Léger, C. Chateau, R. Bini, and L. Ulivi, Ferroelastic phase transition in rutile-type germanium dioxide at high pressure, Phys. Rev. B 58 (6) (1998), pp. R2909–R2912.
- [23] S.P. Mukherjee and S.K. Sharma, Structural studies of gels and gel-glasses in SiO₂-GeO₂ system using vibrational spectroscopy, J. Am. Ceram. Soc. 69 (11) (1986), pp. 806–810.
- [24] P.M. Celliers, G.W. Collins, D.G. Hicks, M. Koenig, E. Henry, A. Benuzzi-Mounaix, D. Batani, D.K. Bradley, L.B. Da Silva, R.J. Wallace, S.J. Moon, J.H. Eggert, K.K.M. Lee, L.R. Benedetti, R. Jeanloz, I. Masclet, N. Dague, B. Marchet, M. Rabec Le Gloahec, Ch. Reverdin, J. Pasley, O. Willi, D. Neely, and C. Danson, *Electronic conduction in shock-compressed water*, Phys. Plasmas 11 (2004), pp. L41–L44.
- [25] N. Goldman, E.J. Reed, I.W. Kuo, L.E. Fried, Ch.J. Mundy, and A. Curioni, Ab initio simulation of equation of state and kinetics of shocked water, J. Chem. Phys. 130 (2009), pp. 14517(1)–124517(6).
- [26] N.C. Colmes, W.J. Nellis, W.B. Graham, and G.E. Walrafen, Spontaneous Raman scattering from shocked water, Phys. Rev. Lett. 55 (1985), pp. 2433–2436.
- [27] R.C. Taylor and G.L. Vidale, The vibrational spectrum of the hydronium ion in hydronium perclorate, J. Am. Chem. Soc. 78 (1956), pp. 5999–6002.
- [28] B. Desbat and P.V. Huong, Spectres i.r. et Raman des sels d'hydronium; $H_3O^+Cl^-, H_3^+Br$ et $H_3O^+SbCl^-$, Spectrochimica Acta, 31A (1975), pp. 1109–1114.
- [29] V.B. Prakapenka, L.S. Dubrovinsky, G. Shen, M.L. Rivers, S.R. Sutton, V. Dimitriev, H.P. Weber, and T. Le Bihan, α-PbO₂-type high-pressure polymorph of GeO₂, Phys. Rev. B 67 (2003), pp. 132101-1–4.
- [30] S.S. Bastanov, Features of solid-phase transformations induced by shock compression, Russian Chem. Rev. 75 (2006), pp. 601–616.